SDO/AIA Observations of a Partially Erupting Prominence

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ABSTRACT

We report an observation of a partially erupting prominence and associated dynamical plasma processes based on observations recorded by the Atmospheric Imaging Assembly (AIA) on board the Solar Dynamics Observatory (SDO). The prominence first goes through a slow rise (SR) phase followed by a fast rise (FR). The slow rise phase started after a couple of small brightenings seen toward the footpoints. At the turning point from SR to FR, the prominence had already become kinked. The prominence shows strong brightening at the central kink location during the start of FR. We interpret this as internal magnetic reconnection occurring at a vertical current sheet forming between the two legs of the erupting prominence (flux-rope). The brightening at the central kink location is seen in all the EUV channels of AIA. The contributions of differential emission at higher temperatures are larger compared to that for typical coronal temperatures supporting a reconnection scenario at the central kink location. The plasma above the brightening location gets ejected as a hot plasmoid-like structure embedded in a CME, and those below drain down in the form of blobs moving towards the Sun’s surface. The unique time resolution of the AIA has allowed all of these eruptive aspects, including SR-to-FR, kinking, central current sheet formation, plasmoid-like eruption, and filament “splitting”, to be observed in a single event, providing strong and comprehensive evidence in favour of the model of partially erupting flux ropes.
1. Introduction

Prominence eruptions are one of the best proxies for Coronal Mass Ejections (CMEs). Therefore, understanding the initiation and evolution of the erupting prominences provides crucial physical understanding of the dynamical processes involved in CME initiation and evolution, with broader implications for space weather and geo-space climate.

Prominence eruptions show a range of observational characteristics from complete eruption to partial eruption to failed or sympathetic eruption. Partial eruption of prominences is a relatively new phenomena which was first described by Gilbert et al. (2000). These partial eruptions were observed in form of downflows from the erupting prominences. Gilbert et al. (2000, 2001) found that majority of the erupting prominences observed with Hα coronagraph located at the Mauna Loa Solar Observatory (MLSO) were associated with downflows. Based on these observations Gilbert et al. (2001) hypothesised that during the process of eruption the prominence breaks into two parts involving an X-type neutral point and internal magnetic reconnection. Depending on the location of internal magnetic reconnection the prominence would show complete, partial or failed eruption (Gilbert et al. 2001).

Gibson and Fan (2006a,b) studied evolution of a flux-rope from the Sun’s surface out to 6 solar radii where a flux-rope with enough twist erupts due to loss of equilibrium. In this simulation a flux-rope undergoes a kink instability leading to the formation of a vertical current sheet between the two legs of the flux-rope (see also Birn et al. 2006). Based on these simulations, Gibson and Fan (2006a,b) demonstrated that multiple internal reconnections (inside the flux-rope) take place at the vertical current sheet. This leads to the bifurcation of the flux-rope. In this process, the part of the flux-rope above the reconnection point escapes as the core of CME and the part of the flux-rope below the reconnection point falls back towards the Sun’s surface. Srivastava et al. (2010) have found observational evidence of kink instability in a highly twisted loop system and the formation of a horizontal current sheet below the double fragmented loop-top.

Based on the above mentioned simulations, Gibson and Fan (2006b) predicted some observables which could be directly compared with observations. Tripathi et al. (2009) studied 6 CMEs events associated with erupting prominences/filaments and concluded that the model of Gibson and Fan (2006a,b) involving partial eruption of flux-rope during the
CME eruption was capable of producing almost all the observed signatures in different
CME events. Thus, the general behaviour described by the model is broadly applicable to
partial eruptions, assisting in interpretation of the CME evolution and its interplanetary
consequences.

Alexander et al. (2006) reported a prominence eruption using the Transition Region
and Coronal Explorer (TRACE; Handy et al. 1999) observations together with Reuven
Ramaty High Energy Solar Spectroscopic Imager (RHESSI; Lin et al. 2002) observations.
They identified two coronal hard X-ray sources, which were located above the filament before
the main activation phase (see also Ji et al. 2003) and right under the apex of the filament
when the filament was strongly kinked. Alexander et al. (2006) attributed this second
brightening to an internal magnetic reconnection occurring at the vertical current sheet
formed under the apex of the kinked filament slightly above the projected crossing point
of its two legs. This event was classified as a failed eruption as there was no associated
CME observed. Liu et al. (2007) reported an observation of partial eruption where the
filament/cavity structure bifurcated into two parts with the upper part ejected as a CME
and the lower part along with the bulk of the filament material remaining on the Sun’s
surface. This event was classified as a case of partially erupting prominence. However, no
clear evidence of reconnection was observed related to this event.

Using multi-wavelength observations recorded by Extreme-ultraviolet Imaging Telescope
(EIT) aboard Solar and Heliospheric Observatory (SOHO), Soft X-ray Telescope (SXT)
aboard Yohkoh, and Hα and white-light coronagraphic observations from the MLSO, Tri-
pathi et al. (2006, 2007) reported observations of a coronal downflow in the wake of a promi-
nence eruption which was bright in EIT and XRT observations. The speed of the downflow
measured using the while-light coronagraphic data was equivalent to the local Alfvén speed.
Furthermore, a cusped shaped structure was observed in SXT images. The downflow started
at the location of X-ray cusp and a base difference images showed an increase in the area of
the dimming region at the cusp. All these observed features were interpreted as signatures
of magnetic reconnection occurring at the location of the cusp. However, due to different
limitations to the observations, the complete evolution of the prominence eruption was not
observed. For instance, the time resolution of EIT or SXT was not enough to show if the
prominence exhibited kinking motions. In addition, neither EIT nor SXT observed the apex
of the cusp because of the limited field of view. On the other hand, Hα and white-light obser-
vations at MLSO showed the apex and location of the bifurcation of the erupting prominence,
but were not available for the early phase evolution, and, since white-light observations are
temperature independent they could not provide information about the thermal properties
of the plasma.
The six events studied by Tripathi et al. (2009) showed many characteristics predicted by the model of Gibson and Fan (2006b). These events separately showed different predicted observables, but none of the events had all the predicted observables together. In addition, due to limited temporal and spatial resolution, some of the predicted characteristics could not be studied, in particular a transition from slow rise (SR) to fast rise (FR), kinking, central current sheet formation, and filament “splitting”.

In the present paper we report an observation of a partially erupting prominence and associated dynamics of the hot plasmoid-like structure using the observations recorded by the Atmospheric Imaging Assembly (AIA) onboard the Solar Dynamics Observatory (SDO). The spatial and temporal resolution of AIA observations and broad temperature coverage is ideal for such observations. In particular, we are able to study the complete evolution and dynamics of an erupting prominence and detect eruptive aspects indicative of a partially-ejected flux rope in the observations. The remainder of the paper is organised as follows. In section 2 we describe the observations followed by the data analysis and results in section 3. We summarise our results and discuss in section 4.

2. Instrumentation and Observations

A small erupting prominence was observed on 12 Sep 2011 by AIA at the East limb of the Sun in a region later named as NOAA AR 11295 on 13 September. AIA provides observations of the Sun in eight different wavelength channels simultaneously with a typical cadence of 12 seconds and pixel size of 0.6 arc sec. The primary contributing ion and the temperature response for each of these channels can vary depending on the features on the Sun’s surface being observed. For more details see O’Dwyer et al. (2011); Del Zanna et al. (2011). In this paper we have investigated all the EUV channels with our main emphasis on 304 Å, 193 Å and 211 Å, and have focussed on the early-phase evolution of the prominence and plasma heating.

Since this is an event occurring on the limb of the Sun, we have also used the observations recorded using SECCHI (Sun Earth Connection and Heliospheric Investigation; Howard et al. 2008) onboard STEREO-B (Solar Terrestrial Relation Observatory), which, because of specific orbit, provides complementary on disk information for the same event. At the time of the observations presented here, the location of STEREO-B spacecraft was such that the features which were seen by SDO at east solar limb, STEREO would see them nearly at the disk centre. Therefore, it provides complementary vantage point for observations.
3. Data Analysis and Results

Figure 1 and the accompanying animation shows a sequence of images taken in 304 Å, passbands of AIA. Note that the images are displayed in negative intensity. The prominence stays relatively quiet until 20:28:44 UT, when two distinct brightenings are seen, marked by two arrows in the top left panel, one near the upper leg of the prominence and the other near the middle. We scrutinised the observations taken in the other bands of AIA and found a similar morphology. The two brightenings (marked by arrows) are strongest in 304 Å channel but also seen mildly in 193 Å and to a lesser extent in 211 Å.

The complete prominence brightens up at around 20:30 UT, about 2 minutes after the first two brightenings, as seen in the 304 images. These images clearly show that the prominence is comprised of many different threads (strands). The brightening in the northern footpoint increases and the prominence goes thorough morphological evolution as revealed by the high time resolution of AIA images. During this time the prominence shows a SR phase. Similar features have previously been reported by Chifor et al. (2006, 2007); Isobe & Tripathi (2006). At 20:40:08 UT (after about 12 mins), when the prominence has attained some height, a large part of the prominence brightens up. At 20:43:20 UT the morphology of the prominence resembles of a kinked flux rope (slinky) with two footpoints attached to the surface. A similar kind of morphological behaviour is seen in the other EUV passbands of AIA. The kinked prominence continues to rise. The prominence seems to go through heating and cooling as manifested by many localised brightenings that again indicate thermally independent strands comprising the prominence. By 20:46:08 UT (after about 3 minutes of the kink), most of the plasma located above the kink location becomes much brighter than the two legs of the prominence that are still connected to the surface (see middle right panel in Fig. 1).

At around 20:46:32 UT (see animation) and 20:46:44 UT (shown by an arrow in bottom left panel in Fig. 1), brightenings took place exactly at the location of the kink. Figure 2 shows near simultaneous observations taken using all six EUV bandpasses of AIA. The brightening shown in the bottom left panel in Fig. 1 is located by an arrow in all the EUV channels of AIA in Fig. 2. This sudden brightening at the kinked location, which is seen in cooler as well hotter channels of AIA, is strongly suggestive of magnetic reconnection occurring at the kink location at the current sheet which may have formed between the two legs of the erupting flux-rope.

Figure 3 displays the height-time plot for the erupting prominence measured using the observations recorded in the 304 Å channel. The plot was obtained by tracking the tip of the erupting prominence. The error bars are due to the standard deviation of three simultaneous measurements of the same features. The plot clearly demonstrates that the prominence goes
through a SR phase up till \(~20:47\) UT before it is accelerated. A linear fit to the data point from the beginning till 20:47 UT yields a speed of about 16 km s\(^{-1}\) and a separate linear fit to the data from 20:47 UT onwards yields a speed of \(~98\) km s\(^{-1}\). The height time plot show a clear 'knee' at around 20:47 UT, which matches with the time of the kinking of the prominence.

We note that the SR to FR change occurs around 20:47 UT, while the kinking becomes evident somewhat earlier around 20:44 UT. Instead of the sharp turning point, there is a knee (cf., Fig. 3) observed that indicates that the kinked prominence threads involved in a bulk reconnection during a span of the time between 20:44-20:47 UT and thereafter prominence rises abruptly. The plasma which is above the location of the kink are ejected outwards and that below drains towards the surface, similar to the observations reported earlier (Tripathi et al. 2006, 2007 see also Gilbert et al. (2001)). The associated CME seen in LASCO/C2 FOV (not shown) first appear at 20:12 UT and disappear quickly within few frames of LASCO/C2 observations. This supports the scenario of reconnection at a central kink location and splitting of the prominence in a partial eruption. Unfortunately, we do not have hard X-ray observations by RHESSI in the similar time domain of the event, therefore, we cannot examine the formation of the HXR at the central kink location as previously observed in various studies (Török et al. 2004; Alexander et al. 2006; Liu and Alexander 2009).

In addition, there are signatures of interaction between the eruptive filament system and the overlying loops, as reflected in form of de-acceleration in the height time diagram shown in Fig. 3. We also note that the overlying fields play important role in prominence eruptions as shown by e.g., Török & Kliem (2005); Fan and Gibson (2007); Kliem (2012).

In order to understand the thermal structure of the erupting prominence and possibly the associated flux rope we have derived emission measure (EM) maps in different temperature bins (see Figure 4) using the six AIA channels. DEMs were calculated using the xrt_dem_iterative2.pro routine (Weber et al. 2004; Golub et al. 2004), modified for use with the AIA passbands. The validation of this method can be found in the appendix of Cheng et al. (2012). The EM was derived at slightly later time than the images shown in Fig. 2 as some of the images were saturated at this time. The emission measure maps were created by calculating a DEM in every pixel using the 6 AIA EUV filters, and then integrating the DEM over the indicated temperature bins in order to get the emission measure in that temperature range. Errors in the DEM calculations are estimated by calculating Monte Carlo runs on the data using intensity values varied normally by the sigma error. This sigma error is estimated by using an empirical formula that approximates Poisson statistics for low count rates, but approaches Gaussian statistics for high count rates (see e.g., Gehrels 1986).
Examining the Monte Carlo runs shows that the errors in the DEMs are higher where the signal in several of the passbands is very low, and where the temperatures are outside the range covered well by the AIA passbands (i.e. <1 MK and > 20 MK).

The EM maps shown in Fig. 4 suggest that the erupting prominence is highly multi-thermal and fairly dense at all temperatures. It should also be noted that the EM is high at the higher temperature bands (e.g. 10-13 MK; the higher red colours at the log EM scale), while comparatively low at typical coronal temperatures (e.g., 1-2 MK; the moderate yellow colour at the log EM scale) at the central kink location. This supports our premises of reconnection generated heating which appears in form of DEM in various temperature bands.

The EM maps also reveal a hot overlying arched structure, which might be representative of an erupting flux rope. Such hot erupting structures have been seen previously by Reeves and Golub (2011); Cheng et al. (2011, 2012, 2013) where they were interpreted as a heated erupting flux rope. However, we do not see any signature of the twist or writhe in this hot overlying flux-rope as previously observed by Liu et al. (2010) and Zhang et al. (2012). Therefore, it may be plausible that this overlying structure is due to the heating at the current sheet between the kinked flux rope and the overlying field (Török et al. 2004). Observations of such a rare heating scenario is possible, most likely due to the narrow EUV passbands of SDO/AIA and availability of the hot channels.

As shown in Figures 1 & 2, the observation reported here is above the limb from the vantage point of SDO. In order to get an on-disk perspective, we have used the observations recorded by STEREO-b. Figure 5 displays images recorded by SECCHI using its 195 Å passband. The spatial and temporal resolution of SECCHI images are not as good as that of AIA. However, the images shown in Figure 5 capture the essential dynamics of the erupting prominence. The arrow in the top right image show the filament in question, which is rising up. The two arrows in the bottom right panel shows that the filament has broken into two parts (shown by two arrows). The bottom part of the filament shows cusp shaped structure. The top of the cusp would correspond to the location of reconnection and enhanced heating, which is captured in the multi-wavelength observations by AIA filters (see Fig. 2). Observations of cusp-shaped structure is one of the most important signatures of magnetic reconnections (see e.g. Forbes and Acton 1996).
4. Summary and Discussion

We have studied complete evolution of a prominence [from early phase evolution to eruption] which shows partial eruption using the excellent observations recorded by AIA on board SDO. The prominence show SR phase during the early phase, followed by a FR. The start of the SR phase coincided with two small brightenings located near the foot points. The prominence shows kinking motion almost at the time of transition from SR to FR. The enhanced brightening was seen at the location of kink, which is seen in all the wavelengths, suggestive of the presence of multi-thermal structures, a premise borne out by the emission measure analysis. The material above the kink location has been observed to be ejected outward (see Fig 3 and associated animation) and the plasma below drains down towards the surface.

To the best of our knowledge, this is the first observation where all the predicted characteristics of partial ejection, including transition from SR-to-FR, kinking, formation of central vertical current sheet, and splitting of filament preceded by magnetic reconnection has been detected in a single event, thanks to the unique time and spatial resolution of AIA. The multi-temperature coverage of the AIA allows us to study the temperature structure of the erupting prominences and study the heating and cooling of plasma during erupting prominences.

These observations provide strong and comprehensive evidence in favour of the model of partially erupting flux ropes. Gibson and Fan (2008) found that the magnetic connectivity, twist, orientation, and topology of the ejected portion of a partially-erupting flux-rope differs significantly from its pre-eruption state. While this idealized simulation cannot capture the details of the observations presented in this paper, the fact that the observations follow the predicted characteristics of that simulation argue that a significant part of the magnetic twist remains behind in the lower portion of a rope that does not erupt; therefore, this region is likely to experience further eruptions. It should also be noted that the observed kink is a configuration of the magnetic flux rope that enables the reconnection but does not necessarily manifest the evolution of the kink instability (Isenberg & Forbes 2007; Török et al. 2010; Liu et al. 2012).

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Fig. 1.— Sequence of AIA images showing the erupting prominence taken in 304 Å passband. Two arrows in the top left image locate the first two brightenings. ‘P’ locates the feature which was tracked to obtain the height-time diagram shown in Fig. The arrow in the bottom left panel located the intense brightening at the location of the kink.
Fig. 2.— Sequence of AIA images recorded with its all the EUV channels showing the brightening at the location of kinked flux-rope.
Fig. 3.— Height-time plot of the erupting prominence measured using the AIA 304 Å observations. This height-time plot was generated by tracking the feature denoted by 'P' in Figure 1.
Fig. 4.— Emission measure maps obtained using six channels of AIA at 20:46:59 UT.
Fig. 5.— Sequence of SECCHI 195 images showing on disk evolution of the erupting filament.